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# Extracting Parton Distribution Functions from data: the Neural Network approach

Andrea Piccione

Edinburgh - Jan, 11th 2006

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#### The NNPDF Collaboration

#### Luigi Del Debbio<sup>1</sup>, Stefano Forte<sup>2</sup>, José I. Latorre<sup>3</sup>, A. P.<sup>4</sup> and Joan Rojo<sup>3</sup>

<sup>1</sup> Particle Physics Theory Group, School of Physics, University of Edinburgh
 <sup>2</sup> Dipartimento di Fisica, Università di Milano and INFN, Sezione di Milano
 <sup>3</sup> Departament d'Estructura i Constituents de la Matèria, Università de Barcelona
 <sup>4</sup> Dipartimento di Fisica Teorica, Università di Torino and INFN, Sezione di Torino

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#### What do we need for the LHC?

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#### What do we need for the LHC?

- SPS: hadronic collider, strong signals and low precision
- ▶ LEP: leptonic collider, *low* signals and high precision
- ► LHC: hadronic collider, *low* signals and high precision

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#### What do we need for the LHC?

- Good reconstruction of final states
- Precise partonic cross-sections calculations
- Accurate description of incoming hadrons

#### How do we describe hadrons?

- QCD describes interactions between quarks and gluons.
   Experimentally we observe only hadrons → Confinement
- Perturbative QCD is not trustable at low energies (~ GeV).
   We can not solve QCD in the non-perturbative region, but on a lattice . . .
- We can extract information on the proton structure from a process with only one initial proton (DIS at HERA).
   Then we can use these as an input for a process where two initial protons are involved (DY at LHC) → Factorization

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### LHC and DIS kinematics



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#### Deep Inelastic Scattering

The cross section

$$\frac{d^2\sigma}{dxdQ^2} = \frac{4\pi\alpha^2}{Q^4} \left[ [1 + (1-y)^2]F_1 + \frac{1-y}{x} (F_2 - 2xF_1) \right]$$

The structure function

$$F_2(x, Q^2) = x \left[ \sum_{q=1}^{n_f} e_q^2 \, \mathcal{C}^q \otimes \boldsymbol{q}_q(x, Q^2) + 2n_f \, \mathcal{C}^g \otimes \boldsymbol{g}(x, Q^2) \right]$$

Parton distribution evlution is described by DGLAP equations

$$Q^2 \frac{d}{dQ^2} q(x, Q^2) = \frac{\alpha_s(Q^2)}{2\pi} (P \otimes q)(x, Q^2)$$

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#### Deep Inelastic Scattering and QCD

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# The problem

- $\blacktriangleright$  For a single quantity  $\rightarrow 1$  sigma error
- For a pair of numbers  $\rightarrow 1$  sigma ellipse
- For a function → We need an "error band" in the space of functions (*i.e.* the probability density P [f] in the space of functions f(x))

Expectation values  $\rightarrow$  Functional integrals

$$\langle \mathcal{F}[f(x)] \rangle = \int \mathcal{D}f \mathcal{F}[f(x)] \mathcal{P}[f(x)]$$

Determine an infinite-dimensional object (a function) from finite set of data points  $\rightarrow$  Mathematically ill-posed problem

Image: A math a math

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# A solution

- 1. Choose a basis of functions, and project the PDFs on it
- 2. Fit coefficients of basis elements by minimizing  $\chi^2$

#### Some possible basis:

[Les Houches Accord PDF sets: Alekhin, Botje, CTEQ, Fermi (GKK), GRV, H1, MRST, ZEUS]

#### Orthogonal Polynomials

[F. J. Yndurain (1978), G. Parisi and N. Sourlas (1979), W. Furmanski and R. Petronzio (1982)]

#### Truncated Moments

[S. Forte and L. Magnea (1999), S. Forte, L. Magnea, A. P. and G. Ridolfi (2001)]

#### Neural Networks

"Any continuous function can be uniformly approximated by a continuous neural network having only one

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#### Some possible basis:

•  $q(x, Q_0^2) = x^{\alpha} (1-x)^{\beta} P(x; \lambda_1, \ldots, \lambda_n)$ 

[Les Houches Accord PDF sets: Alekhin, Botje, CTEQ, Fermi (GKK), GRV, H1, MRST, ZEUS]

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### Still some problems

- Error propagation from data to parameters and from parameters to observables is not trivial
- Theoretical bias due to the choice of a parametrization is difficult to assess (effects can be large if data are not precise or hardly compatible)

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#### The standard approach

MRST: 15 parms. -  $\Delta \chi^2 = 50$  - NC and CC DIS, DY, W-asym, jets

$$xq(x, Q_0^2) = A(1-x)^{\eta}(1+\epsilon x^{0.5}+\gamma x)x^{\delta}, \quad x[\bar{u}-\bar{d}](x, Q_0^2) = A(1-x)^{\eta}(1+\gamma x+\delta x^2)x^{\delta}.$$

$$xg(x, Q_0^2) = A_g(1-x)^{\eta g} (1+\epsilon_g x^{0.5}+\gamma_g x) x^{\delta g} - A_-(1-x)^{\eta} - x^{-\delta} -$$

CTEQ: 20 parms. -  $\Delta \chi^2 = 100$  - NC and CC DIS, DY, W-asym, jets

$$x f(x, Q_0) = A_0 x^{A_1} (1-x)^{A_2} e^{A_3 x} (1+e^{A_4} x)^{A_5}$$

with independent params for combinations  $u_v \equiv u - \bar{u}$ ,  $d_v \equiv d - \bar{d}$ , g, and  $\bar{u} + \bar{d}$ ,  $s = \bar{s} = 0.2 (\bar{u} + \bar{d})$  at  $Q_0$ ; norm. fixed by sum rules

Alekhin: 17 parms. -  $\Delta \chi^2 = 1$  - NC DIS (+ DY)

$$\begin{aligned} xu_V(x, Q_0) &= \frac{2}{N_v^U} x^{a_u} (1-x)^{b_u} (1+\gamma_2^u x); & xu_S(x, Q_0) &= \frac{A_S}{N_S} \eta_u x^{a_s} (1-x)^{b_{SU}} \\ xd_V(x, Q_0) &= \frac{1}{N_d^V} x^{a_d} (1-x)^{b_d}; & xd_S(x, Q_0) &= \frac{A_S}{N^S} x^{a_s} (1-x)^{b_{Sd}}, \\ xs_S(x, Q_0) &= \frac{A_S}{N^S} \eta_S x^{a_s} (1-x)^{(b_{Su}+b_{Sd})/2}; & xG(x, Q_0) &= A_G x^{a_G} (1-x)^{b_G} (1+\gamma_1^G \sqrt{x}+\gamma_2^G x) \end{aligned}$$

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#### The standard approach - Higgs cross section



"Within a given set of PDFs, the deviations of the cross sections from the values obtained with the reference PDF sets can reach the level of 10% at the LHC in the case of the gluon-gluon fusion process for large enough Higgs boson masses,  $M_H \sim 1 TeV$ ".

[A. Djouadi and S. Ferrag, hep-ph/0310209]

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#### The standard approach - HERA-LHC WKS benchmark



"[...] the inclusion of more data from a variety of different experiments moves the central values of the partons in a manner indicating either that the different experimental data are inconsistent with each other, or that the theoretical framework is inadequate for correctly describing the full range of data. To a certain extent both explanations are probably true."

[R. S. Thorne, hep-ph/0511119]

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#### The standard approach - Dependence on $\alpha_s$



"the previous fits with  $\alpha_s(m_Z) =$ 0.118 are adequate for most processes, because the uncertainty associated with  $\alpha_s$  is smaller than the other sources of PDF uncertainty. However,  $\alpha_s$  uncertainty is important for inclusive jet production at relatively small  $p_T$  and Higgs boson production by the  $gg \rightarrow H$  process in the standard model."

[CTEQ, hep-ph/0512167]

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### The NNPDF approach

- ► Determination of the Structure Functions: this is the easiest case, since no evolution is required, but only data fitting. A good application to test the technique → Done
- ► Determination of the Parton Distributions: the real stuff → Working on it ...



#### What are Neural Networks?

Neural networks are a class of algorithms very suitable to fit incomplete or noisy data



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#### Some details on their structure

Building blocks: neurons, *i. e.* input/output units characterized by sigmoid activation

$$\xi_i^{(l)} = g\left(\sum_{j=1}^{n_l-1} \omega_{ij}^{(l-1)} \xi_j^{(l-1)} - \theta_i^{(l)}\right) \quad g(x) = \frac{1}{1 + e^{-x}}$$

- Parameters: weights  $\omega_{ii}^{(l)}$  and thresholds  $\theta_i^{(l)}$ .
- Architecture: multilayer feed-forward NN. Each neuron receives input from neurons in preceding layer and feeds output to neurons in successive layer
- Assumption: smooth function

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Minimization (Training)				

# **Back Propagation**

- 1. Set the parameters randomly.
- 2. Present an input and calculate the output.
- 3. Evaluate  $\chi^2$ .
- 4. Modify the weights to reinforce correct decisions and discourage incorrect ones:

$$\omega_{ij} 
ightarrow \omega_{ij} - \eta rac{\partial \chi^2}{\partial \omega_{ij}}$$

where  $\eta$  is the learning rate.

5. Back to 2, till the stability of  $\chi^2$  is reached

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# Genetic Algorithm

- 1. Set the parameters randomly.
- 2. Make clones of the set of parameters.
- 3. Mutate each clone.
- 4. Evaluate  $\chi^2$  for all the clones.
- 5. Select the clone that has the lowest  $\chi^2$ .
- 6. Back to 2, till the stability of  $\chi^2$  is reached.

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The NNPDF approach				

► Monte Carlo sampling of data (generation of replicas of experimental data) → Faithful error propagation

$$F_i^{(art)(k)} = \left(1 + r_N^{(k)}\sigma_N\right) \left[F_i^{(exp)} + r_i^s\sigma_i^{stat} + \sum_{l=1}^{N_{sys}} r^{l,(k)}\sigma_i^{sys,l}\right]$$

▶ NN training over MC replicas → Unbiased parametrization

Expectation values  $\rightarrow$  Sum over the Nets

$$\left\langle \mathcal{F}\left[F(x,Q^2)\right]\right\rangle = \frac{1}{N_{rep}}\sum_{k=1}^{N_{rep}}\mathcal{F}\left(F^{(net)(k)}(x,Q^2)\right)$$

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#### Details

#### Architecture: 4-5-3-1

- ▶ Inputs: x, log x,  $Q^2$ , log  $Q^2$
- Output:  $F_2(x, Q^2)$

Minimization strategy:

• Back Propagation ( $\sim 10^8$  training cycles):

$$\chi_{\rm diag}^{2\,(k)} = \frac{1}{N_{\rm dat}} \sum_{i=1}^{N_{\rm dat}} \frac{\left(F_i^{(\rm art)(k)} - F_i^{(\rm net)(k)}\right)^2}{\sigma_{i,t}^{(\rm exp)^2}}$$

• Genetic Algorithm ( $\sim 10^4$  generations):

$$\chi^{2(k)} = \frac{1}{N_{\text{dat}}} \sum_{i,j=1}^{N_{\text{dat}}} \left( F_i^{(\text{art})(k)} - F_i^{(\text{net})(k)} \right) \operatorname{cov}_{ij}^{-1} \left( F_j^{(\text{art})(k)} - F_j^{(\text{net})(k)} \right)$$

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Results				
Credits				

- S. Forte, L. Garrido, J. I. Latorre and A. P., "Neural network parametrization of deep-inelastic structure functions," JHEP05 (2002) 062 [arXiv:hep-ph/0204232]
- L. Del Debbio, S. Forte, J. I. Latorre, A. P. and J. Rojo [NNPDF Collaboration], "Unbiased determination of the proton structure function F<sub>2</sub><sup>p</sup> with faithful uncertainty estimation", JHEP03 (2005) 080 [arXiv:hep-ph/0501067]

Source code, driver program and graphical web interface for  $F_2$  plots and numerical computations available @

#### http://sophia.ecm.ub.es/f2neural

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Fit of  $F_2^p(x, Q^2)$ 



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# Same strategy

- Monte Carlo sampling of data
- Parametrization of parton distributions with neural networks
- DGLAP evolution of parton distributions to experimental data scale and training over Monte Carlo replica sample

#### Same strategy, but much more complex!

- Monte Carlo sampling of data
- Parametrization of parton distributions with neural networks
- DGLAP evolution of parton distributions to experimental data scale and training over Monte Carlo replica sample

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#### Examples

Expectation values:

$$\langle \mathcal{F}[g(x)] \rangle = \frac{1}{N_{rep}} \sum_{k=1}^{N_{rep}} \mathcal{F}\left(g^{(net)(k)}(x)\right)$$

$$\sigma_{\mathcal{F}[g(x)]} = \sqrt{\left\langle \mathcal{F}[g(x)]^2 \right\rangle - \left\langle \mathcal{F}[g(x)] \right\rangle^2}$$

 Correlations between pairs of different parton distributions at different points:

$$\langle u(x_1)d(x_2)\rangle = rac{1}{N_{rep}}\sum_{k=1}^{N_{rep}}u^{(net)(k)}(x_1,Q_0^2)d^{(net)(k)}(x_2,Q_0^2)$$

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# PDF Evolution (*very technical!*)

We want Mellin space evolution (numerically efficient):

$$q(N, Q^2) = q(N, Q_0^2) \Gamma\left(N, \alpha_s\left(Q^2\right), \alpha_s\left(Q_0^2\right)\right)$$

We do not want complex neural networks:

$$\Gamma\left(x,\alpha_{s}\left(Q^{2}\right),\alpha_{s}\left(Q^{2}_{0}\right)\right)\equiv\frac{1}{2\pi i}\int_{c-i\infty}^{c+i\infty}dN\,x^{-N}\Gamma\left(N,\alpha_{s}\left(Q^{2}\right),\alpha_{s}\left(Q^{2}_{0}\right)\right)$$

►  $\Gamma(x)$  is a distribution  $\rightarrow$  must be regularized at x = 1:

$$q(x, Q^{2}) = q(x, Q_{0}^{2}) \int_{x}^{1} dy \ \Gamma(y) + \int_{x}^{1} \frac{dy}{y} \Gamma(y) \left( q\left(\frac{x}{y}, Q_{0}^{2}\right) - yq(x, Q_{0}^{2}) \right)$$

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# Details (technical)

- $q_{NS}(x, Q^2) \equiv \frac{1}{6} \left( u + \bar{u} d \bar{d} \right) (x, Q^2)$
- Experimental data: NMC (229 pts) and BCDMS (254 pts)
- Kinematical cuts:  $Q^2 \ge 3 \ GeV^2$ ,  $W^2 \ge 6.25 \ GeV^2$
- ▶ Neural network architecture: 2-5-3-1 (37 params.)
- Strong coupling:  $\alpha_s \left( M_Z^2 \right) = 0.1182$
- Perturbative order: NLO
- ▶ VFN:  $m_c = 1.5 GeV$ ,  $m_b = 4.5 GeV$ ,  $m_t = 175 GeV$
- TMC: F<sub>2</sub> integral evaluated with NN F<sub>2</sub>
- # replica: 100 (should be 1000)

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# Non-Singlet (preliminary)



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# Non-Singlet (preliminary)



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# Non-Singlet (preliminary)



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# Outlook

- Construct full set of NNPDF parton distributions from all available data
- Assess impact of uncertainties of PDFs for relevant observables at LHC
- Make formalism compatible with standard interfaces (LHAPDF, PDFLIB) → NNPDF partons available for use in Monte Carlo generators